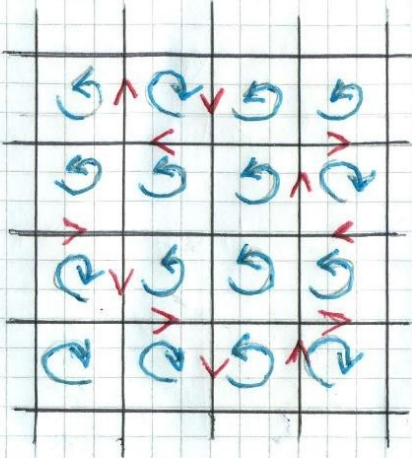


BÁLINT TÓTH
(University of Bristol and Rényi Institute Budapest)

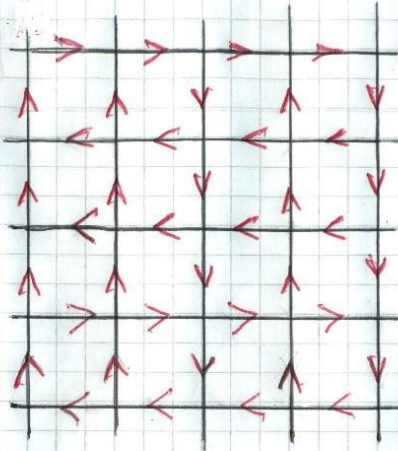
**QUENCHED CLT FOR RANDOM WALK
IN DIVERGENCE-FREE RANDOM DRIFT FIELD**

PDE & Randomness
LMS-Bath, 2021-09-07

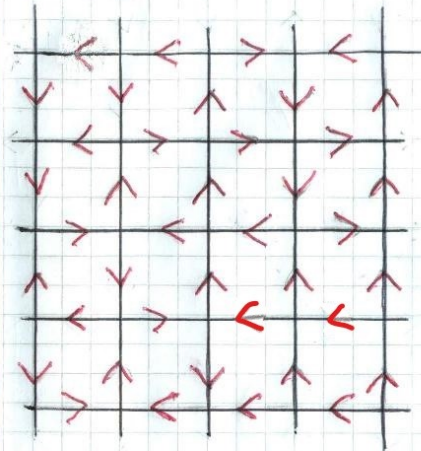
Small cycles
with short
range dependence



"Manhattan"



Six-vertex /
"Square Ice"



$X(t)$: RW, at each step
uniform choice between
the allowed directions

... and also in higher dimension ...

Q: large scale behaviour? diffusive? super diff.?

RW in divergence-free random drift field:

$(\Omega; \pi; \tau_z : \Omega \rightarrow \Omega, z \in \mathbb{Z}^d)$ probab. sp., ergodic \mathbb{Z}^d -action

$\mathcal{E} = \{k \in \mathbb{Z}^d : |k| = 1\}$ possible steps of the rw

$p : \Omega \rightarrow [0, 1]^{\mathcal{E}}$, jump rates of the rw

RWRE: Given $\omega \in \Omega$, $t \mapsto X(t) \in \mathbb{Z}^d$ cont. time Markov chain:

$$\mathbf{P}_\omega (X(t + dt) = x + k | X(t) = x) = p_k(x, \omega) dt = p_k(\tau_x \omega) dt.$$

Separate symmetric and skew-symmetric part of jump rates:

"conductances" :
$$s_k(\omega) := \frac{p_k(\omega) + p_{-k}(\tau_k \omega)}{2} = s_{-k}(\tau_k \omega),$$

"flow" :
$$v_k(\omega) := \frac{p_k(\omega) - p_{-k}(\tau_k \omega)}{2} = -v_{-k}(\tau_k \omega).$$

Assumptions:

$$\sum_{k \in \mathcal{E}} p_{-k}(\tau_k \omega) \equiv \sum_{k \in \mathcal{E}} p_k(\omega) \Leftrightarrow \sum_{k \in \mathcal{E}} v_k(\omega) \equiv 0 \quad \pi\text{-a.s.} \quad (\text{DFR})$$

$$s_k(\omega) \geq s_* > 0, \quad \pi\text{-a.s.} \quad (\text{ELL})$$

$$\int_{\Omega} v_k(\omega) d\pi(\omega) = 0 \quad (\text{NODRIFT})$$

In the examples

$$s_k(\omega) \equiv \frac{1}{2d}, \quad v_k(\omega) = 0, \pm \frac{1}{2d}$$

Analogous diffusion problem: **diffusion in divergence-free random drift field** $t \mapsto X(t) \in \mathbb{R}^d$ with infinitesimal generator

$$L := \frac{1}{2} \nabla \cdot a \nabla + v \cdot \nabla$$

$a(\omega, \cdot) = a(\omega, \cdot)^\dagger : \mathbb{R}^d \rightarrow \mathbb{R}^{d \times d}$; $v(\omega, \cdot) : \mathbb{R}^d \rightarrow \mathbb{R}^d$ are stat+erg,

$$0 < a_* I_{d \times d} \leq a(\cdot, \omega) \leq a^* I_{d \times d} < \infty \quad \pi\text{-a.s.} \quad (\text{ELL})$$

$$\operatorname{div} v(\cdot, \omega) \equiv 0 \quad \pi\text{-a.s.} \quad (\text{DFR})$$

$$\int_{\Omega} v(\cdot, \omega) d\pi(\omega) = 0 \quad (\text{NODRIFT})$$

$$dX(t) = a(\omega, X(t))^{1/2} dB(t) + \left(\frac{1}{2} \nabla a(\omega, X(t)) + v(\omega, X(t)) \right) dt,$$

with $a \equiv I_{d \times d}$: $dX(t) = dB(t) + v(\omega, X(t)) dt$

The (H_{-1})-Condition:

Let

$$C_{kl}(x) := \int_{\Omega} v_k(\omega) v_l(\tau_x \omega) d\pi(\omega), \quad \hat{C}_{kl}(p) := \sum_{x \in \mathbb{Z}^d} e^{\sqrt{-1} p \cdot x} C_{kl}(x)$$

An infrared integrability condition on the correlations:

$$\int_{[-\pi, \pi]^d} \frac{\sum_{k \in \mathcal{E}} \hat{C}_{kk}(p)}{\sum_{i=1}^d (1 - \cos(p_i))} dp < \infty \quad (H_{-1})$$

The natural analytic setting:

$$\Delta : \mathcal{L}^2(\Omega, \pi) \rightarrow \mathcal{L}^2(\Omega, \pi) : \quad \Delta f(\omega) := \sum_{k \in \mathcal{E}} (f(\tau_k \omega) - f(\omega))$$

$$\mathcal{H}_{-1} := \text{Ran}(|\Delta|^{1/2}) = \text{Dom}(|\Delta|^{-1/2}) \subsetneq \mathcal{L}^2(\Omega, \pi)$$

What is the meaning of the (H_{-1}) condition?

$$\mathbf{E}_\omega (dX(t)|X(t) = x) = \underbrace{\sum_{k \in \mathcal{E}} ks_k(\tau_x \omega) dt}_{\psi(\tau_x \omega) \in \mathbb{R}^d} + \underbrace{\sum_{k \in \mathcal{E}} kv_k(\tau_x \omega) dt}_{\varphi(\tau_x \omega) \in \mathbb{R}^d}$$

ψ is a gradient:

$$\psi_i(\omega) = s_{e_i}(\omega) - s_{e_i}(\tau_{-e_i}\omega) \in \mathcal{H}_{-1}$$

The assumption (H_{-1}) is equivalent to: (1) *Analytically*

$$\varphi \in \mathcal{H}_{-1}. \quad (H_{-1})$$

(2) *Probabilistically*: Let $t \mapsto S(t)$ be SSRW independent of ω

$$\lim_{T \rightarrow \infty} T^{-1} \int_{\Omega} \mathbf{E} \left(\left| \int_0^T \varphi(\tau_{S(t)}\omega) dt \right|^2 \right) d\pi(\omega) < \infty. \quad (H_{-1})$$

First consequences of the assumptions?

- (DFR) \Rightarrow The *environment process*:

$$t \mapsto \eta_t := \tau_{X(t)}\omega$$

is stationary and ergodic Markov process in (Ω, π) .

- (H_{-1}) \Rightarrow (NODRIFT) \Rightarrow zero overall drift of the walk:

$$\mathbf{E}(X(t)) := \int_{\Omega} \mathbf{E}_{\omega}(X(t)) d\pi(\omega) = 0.$$

- Thus, (DFR) & (NODRIFT) \Rightarrow (by the ergodic thm) LLN:

$$t^{-1}X(t) \rightarrow 0, \quad \text{a.s.}$$

- (ELL) & (H_{-1}) \Rightarrow diffusive bounds (not straightforward):

$$0 < \underline{\lim}_{t \rightarrow \infty} t^{-1} \mathbf{E}(|X(t)|^2) \leq \overline{\lim}_{t \rightarrow \infty} t^{-1} \mathbf{E}(|X(t)|^2) < \infty.$$

Question left open: CLT?

Why care?

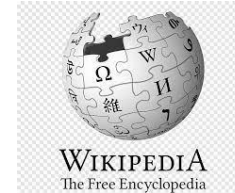
- (1) Genuinely **more general** than RW among random conductances. Even the $s \equiv 1$ case is interesting and non-trivial.
- (2) **Physical motivation**: Diffusion of suspended particles in *incompressible turbulent flow*, in stationary regime.
- (3) Assuming a.c. stationary measure $\rho(\omega)d\pi(\omega)$ for $t \mapsto \eta(t) \in \Omega$, a random change of time leads to RW in div-free drift field.

Why different?

Essentially (but not only) due to **non-self-adjointness**.

History

From **Lucretius Carus (cca 60 BC)** to
G Papanicolaou, SRS Varadhan (1981) (more later)



From “Brownian Motion” at

History[[edit source](#)]

The Roman philosopher-poet [Lucretius](#)' scientific poem "[On the Nature of Things](#)" (c. 60 BC) has a remarkable description of the motion of [dust](#) particles in verses 113–140 from Book II. He uses this as a proof of the existence of atoms:

.....

Although the mingling motion of dust particles is caused largely by air currents, the glittering, tumbling motion of small dust particles is caused chiefly by true Brownian dynamics; Lucretius "perfectly describes and explains the Brownian movement by a wrong example".^[8]

I don't think it is a "wrong example" of BM.
It shows universality at very different length-scales.

Titus Lucretius Carus: De Rerum Natura, Liber Secundus (cca 60 BC)

sic a principiis ascendit motus et exit // paulatim nostros ad
sensus, ut moveantur // illa quoque, in solis quae lumine cernere
quimus // nec quibus id faciant plagis apparet aperte.

Thus motion ascends from the primevals on, // And stage by
stage emerges to our sense, // Until those objects also move
which we // Can mark in sunbeams, though it not appears //
What blows do urge them.

(Transl. William Ellery Leonard, 1916)

Proposition ("Helmholtz's theorem"). Let $v_k(x, \omega) := v_k(\tau_x \omega)$ be an ergodic, \mathcal{L}^2 , divergence-free flow on \mathbb{Z}^d (as before).

(i) $\exists!$ tensor-field $x \mapsto h(x, \cdot) \in \mathcal{L}^2(\Omega \rightarrow \mathbb{R}^{\mathcal{E} \times \mathcal{E}}, \pi)$ such that

cocycle $h_{k,l}(y, \omega) - h_{k,l}(x, \omega) = h_{k,l}(y - x, \tau_x \omega) - h_{k,l}(0, \tau_x \omega),$

tensor $h_{l,k}(x, \omega) = h_{-k,l}(x + k, \omega) = h_{k,-l}(x + l, \omega) = -h_{k,l}(x, \omega),$

v = roth h $v_k(x, \omega) = \sum_{l \in \mathcal{E}} h_{k,l}(x, \omega).$

($x \mapsto h(x)$ is the stream tensor.)

(ii) (H_{-1}) is valid (for v) if and only if for some $h_{k,l} \in \mathcal{L}^2(\Omega, \pi)$

$$h_{k,l}(x, \omega) = h_{k,l}(\tau_x, \omega).$$

[Comment: uniqueness in (i) = Coulomb gauge]

Theorem. Assume **(DFR)**, **(ELL)**, **(H₋₁)**. The non-degenerate covariance matrix $(\sigma^2)_{ij} := \lim_{T \rightarrow \infty} T^{-1} \mathbf{E} \left(X_i(T) X_j(T) \right)$ exists.

For any $m \in \mathbb{N}$, $t_1, \dots, t_m \in \mathbb{R}_+$ and $f : \mathbb{R}^{md} \rightarrow \mathbb{R}$ cont & bdd

(i) [G Kozma, BT (*Ann. Probab.* 2017)]

CLT in probability w.r.t. the environment – *semi-quenched*:

$$\mathbf{E}_\omega \left(f(\dots, T^{-1/2} X(Tt_j), \dots) \right) \xrightarrow{\pi\text{-prob}} \mathbf{E} \left(f(\dots, W_\sigma(t_j), \dots) \right).$$

(ii) [BT (*Ann. Probab.* 2018)]

CLT almost surely w.r.t. the environment – *quenched*:

If the marginally stronger integrability condition

$$\int_{\Omega} |h_{k,l}(\omega)|^{2+\varepsilon} d\pi(\omega) < \infty \quad (\mathbf{H-1-turbo})$$

holds, then

$$\mathbf{E}_\omega \left(f(\dots, T^{-1/2} X(Tt_j), \dots) \right) \xrightarrow{\pi\text{-a.s.}} \mathbf{E} \left(f(\dots, W_\sigma(t_j), \dots) \right).$$

Historical comments (sketchy, far from complete):

[SM Kozlov (1979)], [G Papanicolaou, SRS Varadhan (1981)]:

$v \equiv 0$, self-adjoint, diffusion, initiation of the problem

[H Osada (1983)], [SM Kozlov (1985)]:

$h \in \mathcal{L}^\infty$, \mathbf{O} : quenched diffusion; \mathbf{K} : semi-quenched walk

[K Oelschläger (1988)], [M Avellaneda, A Majda (1991)],

[A Fannjiang, G Papanicolaou (1996)]:

$s = \text{const.}$, $h \in \mathcal{L}^2$, semi-quenched, diffusion, some restrictions

[A Fannjiang, T Komorowski (1997)]:

$s = \text{const.}$, $h \in \mathcal{L}^{d+\varepsilon}$, quenched diffusion.

[T Komorowski, S Olla (2003)], [J-D Deuschel, H Kösters (2008)]:

$h \in \mathcal{L}^\infty$, \mathbf{KO} : semi-quenched walk, \mathbf{DK} : quenched walk

[T Komorowski, C Landim, S Olla (2012)]:

$h \in \mathcal{L}^d$, semi-quenched walk, \dagger diffusion with s, h Gaussian

Elements of proof:

- **Quenched tightness:** Nash moment bound **extended** from $h \in \mathcal{L}^\infty(\Omega, \pi)$ to $h \in \mathcal{L}^{2+\varepsilon}(\Omega, \pi)$:

$$\overline{\lim}_{t \rightarrow \infty} t^{-1/2} \mathbf{E}_\omega (|X(t)|) \leq C < \infty, \quad \pi\text{-a.s.}$$

- **Harmonic coordinates:** Kozlov, Osada . . . **extended** from $h \in \mathcal{L}^\infty(\Omega, \pi)$ to $h \in \mathcal{L}^2(\Omega, \pi)$:

Given $\varphi \in \mathcal{H}_{-1}$,

find $\theta \in \mathcal{L}^2(\Omega \rightarrow \mathbb{R}^\mathcal{E}, \pi)$, s.t. $\theta_k(\omega) + \theta_{-k}(\tau_k \omega) = 0$

$$\sum_{k \in \mathcal{E}} p_k(\omega) \theta_k(\omega) = \varphi,$$

$$\theta_k(\omega) + \theta_l(\tau_k \omega) = \theta_l(\omega) + \theta_k(\tau_l \omega)$$

Let $x \mapsto \Theta(x, \omega)$ be the \mathbb{Z}^d -cocycle with $\text{grad } \Theta(x, \omega) = \theta(\tau_x \omega)$

Elements of proof (ctd):

- **"Relaxed Sector Condition" (RSC):** Prove that

$|\Delta|^{-1/2} (L - L^*) |\Delta|^{-1/2}$ is **skew-self-adjoint** over $\mathcal{L}^2(\Omega, \pi)$.

- **Quenched martingale CLT** for

$$t \mapsto X(t) - \Theta(X(t), \omega)$$

(given the harmonic coordinates, this comes for free)

- **Control the corrector:**

$$\lim_{t \rightarrow \infty} \mathbf{P}_\omega \left(t^{-1/2} |\Theta(X(t), \omega)| > \delta \right) = 0, \quad \pi\text{- a.s.}$$

relying on quenched tightness and [Zygmund (1951)]'s **unrestricted a.s. ergodic theorem** (over \mathbb{Z}^d).

RSC – functional analytic details: (assume $s \equiv 1$, a.s.)

Some **operators** on the Hilbert space $\mathcal{L}^2(\Omega, \pi)$:

$$\mathcal{L}^2(\Omega, \pi)\text{-grad} : \quad \nabla_k f(\omega) := f(\tau_k \omega) - f(\omega), \quad \nabla_k^* = \nabla_{-k}$$

$$\mathcal{L}^2(\Omega, \pi)\text{-Lapl} : \quad \Delta f(\omega) := \sum_{k \in \mathcal{E}} (f(\tau_k \omega) - f(\omega)). \quad \Delta^* = \Delta \leq 0$$

$$\text{multipl. ops.} : \quad M_k f(\omega) := v_k(\omega) f(\omega), \quad M_k^* = M_k$$

and a commutation relation – due to (**DFR**):

$$\sum_{k \in \mathcal{E}} M_k \nabla_k + \sum_{k \in \mathcal{E}} \nabla_{-k} M_k = 0$$

The **infinitesimal generator** of the environment process:

$$L = \frac{1}{2} \Delta + \sum_{k \in \mathcal{E}} M_k \nabla_k = \frac{1}{2} \Delta - \sum_{k \in \mathcal{E}} \nabla_{-k} M_k =: -S + A$$

Formal solution of the eq. for gradient of harmonic coordinates:

$$\theta_k = |\Delta|^{-1/2} \nabla_k \left(I + \underbrace{|\Delta|^{-1/2} A |\Delta|^{-1/2}}_B \right)^{-1} |\Delta|^{-1/2} \varphi$$

Note:

$$\left\| |\Delta|^{-1/2} \nabla_k \right\| \leq 1, \quad \left\| |\Delta|^{-1/2} \varphi \right\| < \infty, \quad B \text{ unbounded.}$$

Target: Prove that

$$B := \sum_{k \in \mathcal{E}} (|\Delta|^{-1/2} \nabla_{-k}) M_k |\Delta|^{-1/2}$$

$$B^* := |\Delta|^{-1/2} \sum_{k \in \mathcal{E}} M_k (\nabla_k |\Delta|^{-1/2}),$$

defined on appropriately chosen cores $\mathcal{B}, \mathcal{B}^* \subset \mathcal{L}^2(\Omega, \pi)$, make sense as *unbounded, skew-self-adjoint* operator.

Checking **von Neumann's criterion** $\text{Ker}(I \pm B^*) = \{0\}$:
boiles down to proving that there is no (non-trivial) solution to

$$\sum_{k \in \mathcal{E}} p_k(\omega) \theta_k(\omega) = 0,$$

$$\theta_k(\omega) + \theta_{-k}(\tau_k \omega) = 0$$

$$\theta_k(\omega) + \theta_l(\tau_k \omega) = \theta_l(\omega) + \theta_k(\tau_l \omega)$$

Gometrically: No non-trivial harmonic cocycles!

By-product: Uniqueness of the harmonic coordinates.

When (H_{-1}) fails: expect superdiffusive behaviour.

Diffusion in the curl of 2-dim Gaussian Free Field:

$$v : \mathbb{R}^2 \mapsto \mathbb{R}^2, \quad v = \text{curl}(U * GFF), \quad dX(t) = dB(t) + v(X(t))dt.$$

Theorem. [BT, B Valkó (*Journ. Stat. Phys.* 2012)]

$$d = 2 : t \log \log t \ll \mathbf{E}(|X(t)|^2) \ll t \log t, \quad (\text{expect } \asymp t\sqrt{\log t})$$

annealed, in the sense of Laplace transform (modulo Tauberian inversion).

Recent development:

G. Cannizzaro & L. Haunschmid-Sibitz & F. Toninelli (2021):

$$\mathbf{E}(|X(t)|^2) \asymp t\sqrt{\log t}$$

Go to **Fabio**'s lecture on Friday.

RW on random Manhattan lattice:

Theorem. [S Ledger, BT, B Valkó (ECP 2018)]

$$d = 2 : t^{5/4} \ll \mathbf{E}(|X(t)|^2) \ll t^{3/2}, \quad (\text{expect } \asymp t^{4/3}, \text{WIP})$$

$$d = 3 : t \log \log t \ll \mathbf{E}(|X(t)|^2) \ll t \log t, \quad (\text{expect } \asymp t\sqrt{\log t}, \text{WIP})$$

annealed, in the sense of Laplace transform (modulo Tauberian inversion).